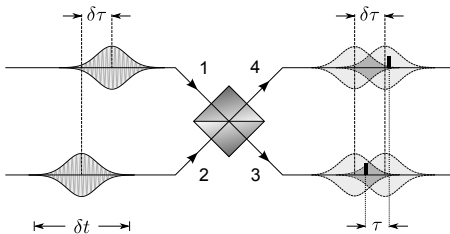


Two-Photon Interference

Markus Perner

Technische Universität München

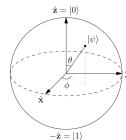
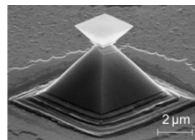
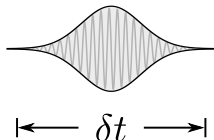
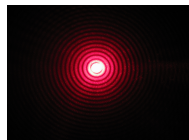
08. Mai 2013



Motivation

Why is two-photon interference so important?

- Show quantum nature of light
- Measurement of the length of a photon
- Characterization of single-photon sources
- Linear quantum information processing (QIP)



Outline

- 1 Theory
- 2 Hong-Ou-Mandel Interference [Hong, Ou & Mandel (1987)]
- 3 Quantum Beat of Two Single Photons [Legero et al. (2004)]
- 4 Summary

Outline

1 Theory

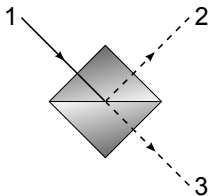
2 Hong-Ou-Mandel Interference [Hong, Ou & Mandel (1987)]

3 Quantum Beat of Two Single Photons [Legero et al. (2004)]

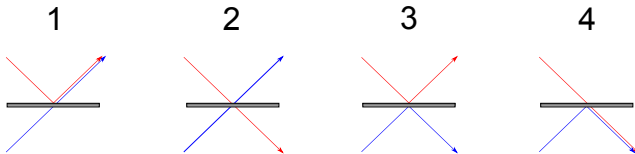
4 Summary

Physical Description

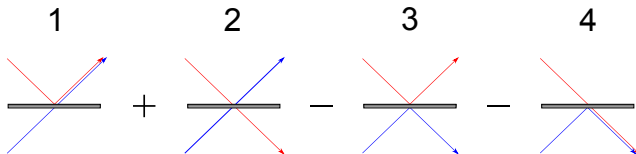
When a photon enters a beam splitter (BS), there are two possibilities: it will either be transmitted or reflected



Consider two photons, one in each input mode of a 50:50 BS. There are four possibilities for the photons to behave:



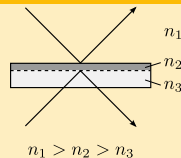
The state of the system after interference is given by a superposition of all possibilities for the photons to pass through the BS:



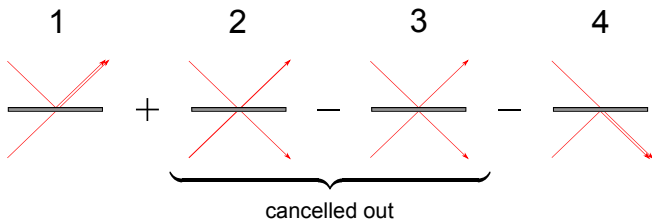
Where does the minus sign come from?

Reflection off bottom side $\Rightarrow$ relative phase shift of π
(i.e. reflection off the higher index medium)

Reflection off top side $\Rightarrow$ no phase shift
(i.e. reflection off the lower index medium).



Now let's assume that the two photons are **identical** in their physical properties (i.e., polarization, spatio-temporal mode structure, and frequency):



The photons will always exit the same (but random) output port!

Mathematical Description

- Consider two identical photons input ports A and B of 50:50 BS
- Quantum state given by applying **photon-creation operators** $\hat{a}_{A,B}^\dagger$ on vacuum state $|0\rangle$

$$|\Psi_i\rangle = |1_A 1_B\rangle = \hat{a}_A^\dagger \hat{a}_B^\dagger |0\rangle$$

- Detection of photon in output port C or D evaluated by applying **photon-annihilation operators** $\hat{a}_{C,D}$ to $|\Psi_i\rangle$
- Effect of 50:50 BS described by **unitary transformation**

$$\begin{pmatrix} \hat{a}_A \\ \hat{a}_B \end{pmatrix} \rightarrow \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} \hat{a}_C \\ \hat{a}_D \end{pmatrix}$$

Initial state in terms of photons created in output ports:

$$\begin{aligned}
 |1_A 1_B\rangle &= \hat{a}_A^\dagger \hat{a}_B^\dagger |0\rangle \\
 &= \frac{1}{2} (\hat{a}_C^\dagger + \hat{a}_D^\dagger) (\hat{a}_C^\dagger - \hat{a}_D^\dagger) |0\rangle \\
 &= \frac{1}{2} (\hat{a}_C^{\dagger 2} - \hat{a}_D^{\dagger 2}) |0\rangle \\
 &= \frac{1}{\sqrt{2}} (|2_C 0_D\rangle - |0_C 2_D\rangle)
 \end{aligned}$$

Remark: Since the operators $\hat{a}_{C,D}^\dagger$ act on different output ports their commutator vanishes: $[\hat{a}_C^\dagger, \hat{a}_D^\dagger] = 0$

The photons will always exit the same (but random) output port!

Outline

1 Theory

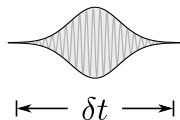
2 Hong-Ou-Mandel Interference [Hong, Ou & Mandel (1987)]

3 Quantum Beat of Two Single Photons [Legero et al. (2004)]

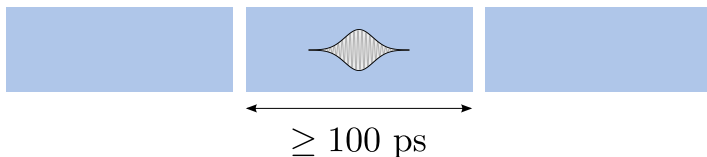
4 Summary

Hong-Ou-Mandel Interference

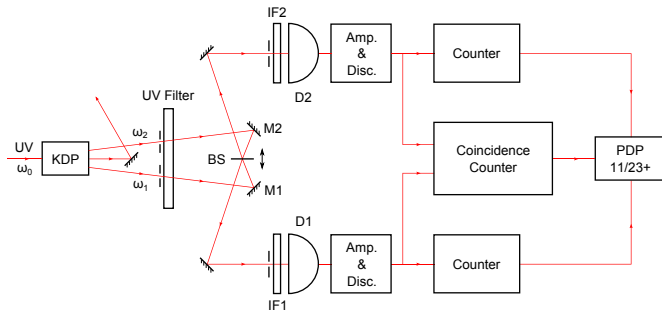
Aim: Measurement of photon wave packet length δt



Problem: Time resolution $\gg$ wave packet length



Experimental setup



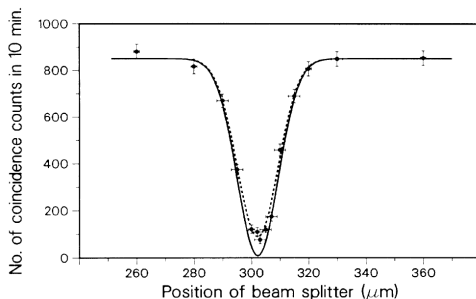
Photon source: Parametric down-conversion Type I ($\omega_0 = \omega_1 + \omega_2$)

Indistinguishability of photons:

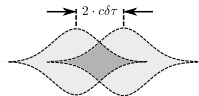
- Photon source ensures **same polarization**
- Pass bands of interference filters (IF) ensure (at most) **same frequency**
- BS is displaced (by $\pm c\delta\tau$) from symmetry position to ensure **same time of arrival**

Measurement of simultaneous detections N depending on BS displacement $\delta\tau$

Hong-Ou-Mandel-Dip



- Overlap controlled by BS displacement:
no coincidences for perfect overlap
- Photon length from FWHM = $16 \mu\text{m} = c \cdot 50 \text{ fs}$
- Width of dip connected to
IF-bandwidth due to Fourier-limited
photons



Outline

1 Theory

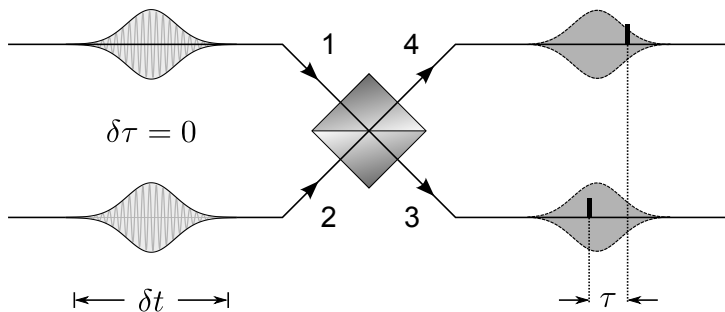
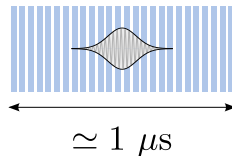
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4 Summary

Quantum Beat of Two Single Photons

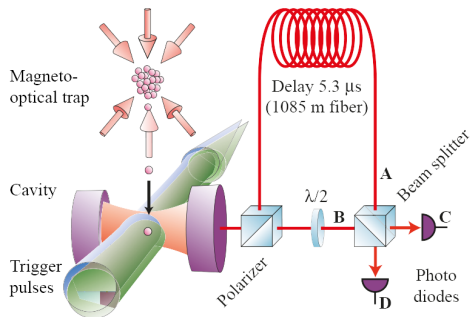
Now: Time resolution $\ll$ wave packet length
 Photons impinge simultaneous on BS
 Detection-time delay τ measured



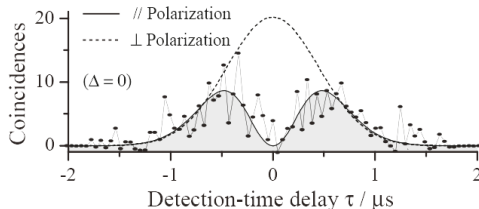
Experimental Setup

Photon source: Atom-cavity system

- Atom-cavity system emits **unpolarized single** photons with **well determined frequency**
- Polarizing beam splitter (PBS) **randomly** directs them along two paths
- Photon travelling along path *A* gets **delayed** through fiber
- Subsequent photon travelling along path *B* impinges on the BS **simultaneously** with delayed photon from path *A*
- A HWP along path *B* guarantees that both photons have the **same polarization**



Results 1

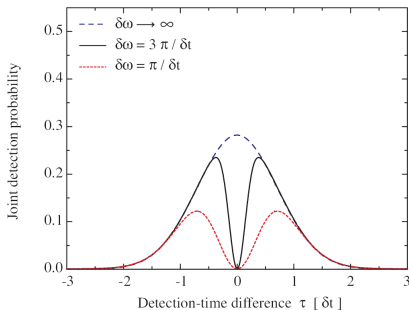
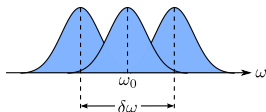


- Indistinguishable photons for parallel polarization: Correlation drops to minimum with width of 460 ns
- Width corresponds to inhomogeneous broadening of photon spectrum of $\delta\omega/2\pi = 690$ kHz
- In general depth and width of minimum indicate initial purity of photon state, i.e. broadening of photon spectrum, frequency and emission-time jitter

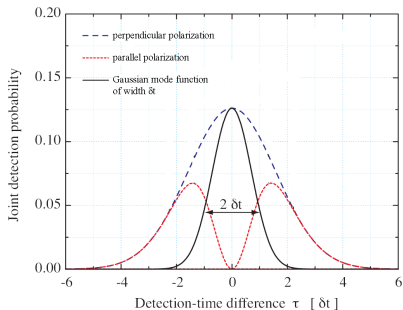
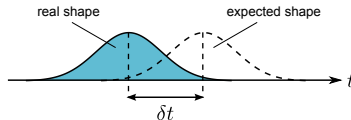
Jitters

No perfect single-photon source $\implies$ Consider more realistic scenario, in which stream of Fourier-limited photons shows a **jitter** in its parameters:

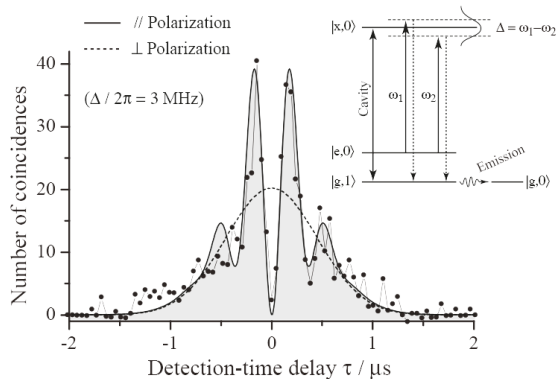
■ Frequency jitter:



■ Emission-time jitter:



Results 2



- Introducing a frequency difference of $\Delta\omega/2\pi = 3 \text{ MHz}$: Quantum beat visible
- No coincidences for simultaneous detection
- Fringes with a visibility of almost 100% (classical: 50%)

Theoretical Prediction of Quantum Beat

- Initial state: $|\Psi_i\rangle = |1_A 1_B\rangle$
- Detection of photon at time t_0 in output C :

$$|\Psi_C(t_0)\rangle = \hat{a}_C |1_A 1_B\rangle = \frac{1}{\sqrt{2}}(|1_A 0_B\rangle + |0_A 1_B\rangle)$$

- Introduce frequency difference $\Delta = \omega_2 - \omega_1$. After time τ new state reads:

$$|\Psi_C(t_0 + \tau)\rangle = \frac{1}{\sqrt{2}}(e^{i\omega_1\tau} |1_A 0_B\rangle + e^{i\omega_2\tau} |0_A 1_B\rangle)$$

- Introduce global phase $|\Psi_C\rangle \rightarrow e^{-i\omega_1\tau} |\Psi_C\rangle$:

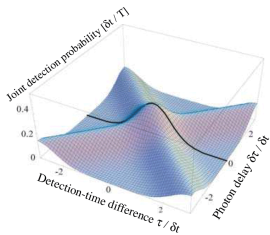
$$\Rightarrow |\Psi_C(t_0 + \tau)\rangle = \frac{1}{\sqrt{2}}(|1_A 0_B\rangle + e^{i\Delta\cdot\tau} |0_A 1_B\rangle)$$

- Probability of detecting the second photon with either C or D :

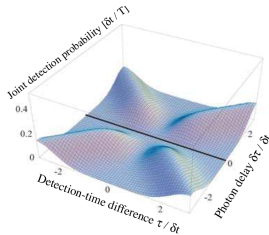
$$P_{CC} = \langle \Psi_C | \hat{a}_C^\dagger \hat{a}_C | \Psi_C \rangle = \frac{1}{2} [1 + \cos(\Delta \cdot \tau)]$$

$$P_{CD} = \langle \Psi_C | \hat{a}_D^\dagger \hat{a}_C | \Psi_C \rangle = \frac{1}{2} [1 - \cos(\Delta \cdot \tau)]$$

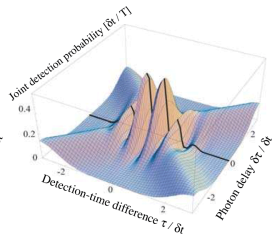
A frequency difference results in a quantum-beat signal



perpendicular polarization:
no interference



parallel polarization:
interference



interference of photons with
different frequencies

no simultaneous detections
($\tau = 0$) for $\delta \tau = 0$

even for distinguishable
photons no simultaneous
detections ($\tau = 0$)

oscillation of coincidence
probability for $\delta \tau \simeq 0$

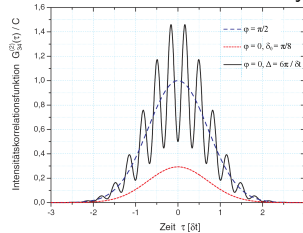
Comparison to HOM-Experiment:

No time resolved measurement $\implies$ Integrate over detection-time difference τ :

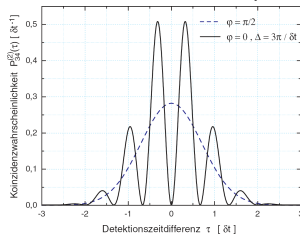
No difference between cases 1 and 3

Visibility

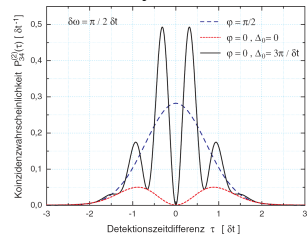
Classical fields: 50% visibility



QM two-photon interference: 100% visibility



Fourier-limited photons



Inhomogeneous broadened photons

Outline

1 Theory

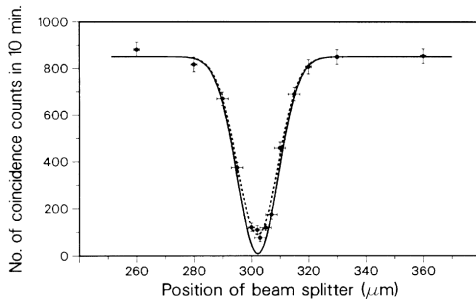
2 Hong-Ou-Mandel Interference [Hong, Ou & Mandel (1987)]

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4 Summary

Summary

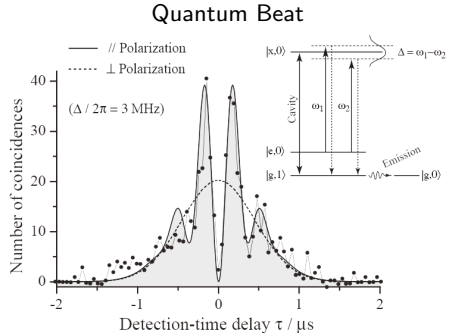
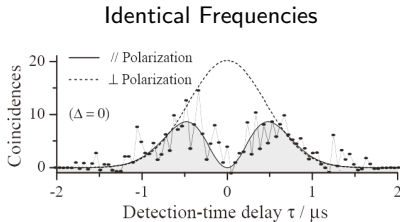
Hong-Ou-Mandel Interference Photon source: PDC



- No time-resolved measurement possible
- Width of dip connected to length of photon wave packet

Two-Photon Interference

Photon source: Single Photon Emitter



- Time-resolved measurement possible
- Photon source can be characterized from width and depth of dip (i.e. broadening of photon spectrum, frequency and emission-time jitter)
- Frequency difference between photons induces quantum-beat with visibility of 100%

Thank you for your attention
and I look forward for your questions on this topic!

References



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